Low-loss light transmission in a rectangularshaped hybrid metal trench at 1550 nm

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Abstract: A hybrid plasmonic waveguide consisting of a high-index dielectric core embedded inside a rectangular-shaped metallic trench is proposed and its guiding properties are investigated at the wavelength of 1550 nm. Numerical simulations based on the finite element method have demonstrated that the introduced dielectric core could greatly reduce the modal loss of the metal trench while maintaining strong confinement of light. The effects of dielectric core size, material of the cladding and the dielectric core on the modal properties have been systematically investigated. The proposed hybrid plasmonic structure can be realized employing fabrication techniques of the traditional metal trench waveguides and could be leveraged as important elements for highly-integrated photonic circuits.

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1. Introduction

The goal of manipulating light at the subwavelength scale has motivated extensive research into a number of waveguiding structures that could support light propagation well beyond the fundamental diffraction limit [1]. Surface plasmon polariton (SPP) based waveguides are promising candidates for subwavelength light guiding due to their capabilities of providing truly nanoscale confinement along two dimensions. Owing to their fancinating optical properties, they have been considered as important building blocks for novel miniaturized photonic components and circuits [2].

Up to date, a number of SPP waveguiding structures have been proposed and investigated, which include metallic nanowires [3], metal stripes embedded in homogeneous dielectrics [4], metal slots [5] or channels [6, 7] cut into metal substrate as well as triangular metal wedges [8]. However, for most of these conventional SPP waveguides, the tradeoff between the mode confinement and transmission loss could not be well balanced. Metal slot waveguide [5] could offer nanoscale confinement of the lightwave but its modal loss is relatively high due to the

significantly increased field strength in the metallic structure. On the other hand, metallic stripe structures could support the propagation of ultra-low-loss plasmonic modes. However, their confinement is weak and the mode size is typically larger than the wavelength scale [4].

Recently, a novel hybrid plasmonic waveguide that leverages the coupling between dielectric and plasmonic modes has demonstrated better capabilities to compromise the propagation loss and field confinement [9]. Through introducing an additional low-index layer between the high-index structure and the metal waveguide, the strong coupling between different optical modes has resulted in a novel hybridized mode featuring strong field localization near the low-index gap with low modal loss. Its nice optical performance has enabled the realization of plasmon nanolasers [10], and also facilitated the operation of numerous ultra-compact passive plasmonic devices [11-15]. Furthermore, a number of modified hybrid structures have also been intensively studied, which include metal-dielectricslabs [16–20], dielectric-loaded type [21–23], dielectric-rib waveguide [24], coaxial structure [25, 26], nanowire array [27], metal-insulator-metal type [15, 28–31], and wedge [32] or ridge [33] configurations.

Here in this paper, we would like to leverage the guiding capability of a rectangular metallic trench [7]. By combining high-index dielectric with a metal trench structure, a novel hybrid plasmonic waveguide featuring both low propagation loss and strong optical confinement can be realized. In the following section, we will carry out detailed investigations on the properties of the fundamental hybrid plasmonic modes against the changes of key geometric parameters at the telecom wavelength.

2. Geometries and modal properties of the hybrid metal trench waveguide

Figure 1 schematically shows the geometry of the proposed hybrid trench waveguide, where a rectangular high-index dielectric core is embedded inside the metal trench, with thin lowindex gaps between the core and the sidewalls of the metal trench. The materials of the lowindex gaps and the cladding are assumed to be the same. The depth of the metal trench and the height of the dielectric core are denoted as T. The width of the dielectric core is W_d . The gap distance between the dielectric core and the sidewall is W_s . The low-index dielectric layer has a thickness T. The refractive indices of the metal trench, dielectric core and the cladding are n_m , n_d and n_c , respectively. The modal properties of the hybrid waveguides are investigated at $\lambda = 1550$ nm by the finite-element method (FEM) using COMSOLTM. Scattering boundary conditions along with extremely fine mesh are applied near the gap regions to ensure accurate solutions. In the simulations, the metal trench is assumed to be silver (Ag) and its refractive index is $n_m = 0.1453 + 11.3587i$ [27]. The high-index core is first assumed as silicon and has a refractive index of 3.5. In the following studies, the depth of the metal trench (i.e. height of the dielectric core) is set at T = 400 nm without specific notifications.



Fig. 1. Schematic illustration of the studied hybrid trench waveguide.

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We firstly look into the propagation property of the proposed hybrid trench structure, where the results of the conventional metal trench without the high-index core and the proposed hybrid structure are both shown in Fig. 2. The propagation length is calculated by

$$L = \lambda / [4\pi \operatorname{Im}(N_{eff})] \tag{1}$$

Here, $\text{Im}(N_{eff})$ is the imaginary part of the modal effective index. It can be seen from the figure that, the propagation lengths for both cases would reduce monotonically with the increase of the refractive index of the cladding. Figure 2 also illustrates that, compared to the conventional metal trench, the proposed hybrid trench has lower modal loss, and its propagation length can be more than 2 times larger than the conventional one. For instance, as show in the figure, for $W_s = 25$ nm and $W_s = 50$ nm, when $n_c = 1$, the propagation length are 15 µm and 25 µm for conventional metal trench waveguide, and are 55 µm and 60 µm for our proposed hybrid structure. On the other hand, when $n_c = 3$, the propagation lengths of the conventional metal trench and hybrid trench are 5 µm and 20 µm for $W_s = 25$ nm, and 8 µm and 22 µm for $W_s = 50$ nm, respectively.



Fig. 2. Dependence of the propagation lengths of the conventional metal trench and hybrid trench waveguides on the refractive index of the cladding layer.

To explain the difference between the modal losses of the two structures, we further show the profiles of the two plasmonic modes. In Fig. 3, the energy density distributions of the conventional metal trench structure and the proposed hybrid trench waveguide are depicted. It is clearly illustrated that, for the proposed hybrid waveguide, a large portion of the modal energy can be confined in the dielectric core, thereby resulting in lower propagation loss. However, for the conventional metal trench waveguide, due to the strong confinement in the metal, its modal loss is much larger than the hybrid case.



Fig. 3. Normalized energy density distributions of (a) conventional metal trench waveguide; (b) hybrid trench waveguide.

We then investigate the propagation properties of the hybrid trench mode for structures with dielectric cores of different size. Here the transmission loss can be calculated by the following equation [10], and the calculated results of the propagation length and the transmission loss are shown in Fig. 4(a) and 4(b), respectively.



Fig. 4. Propagation properties for waveguides with different dielectric core sizes and claddings: (a) propagation length and (b) transmission loss.

It can be seen from Fig. 4 that, for small W_d and small gap width, increasing the refractive index of the cladding would lead to the monotonic decrease of the propagation length and increase in the transmission loss. However, when both the dielectric core and the gap width are relatively large (e.g. $W_d = 300$ nm, $W_s = 50$ nm or $W_d = 200$ nm, $W_s = 50$ nm), the changes of the propagation properties demonstrate non-monotonic trends.

In order to explain the above phenomenon more clearly, we need to elaborate the basic working principle of our waveguide. The coupling between the dielectric mode and plasmonic mode plays an important role in determining the modal characteristics. Owing to the presence of the high-index dielectric core, the energy that originally concentrated in two independent slit waveguide gets strongly coupled in the dielectric core. Two physical parameters including the refractive index of cladding material and the size of the high-index core (core width) will significantly influence the coupling efficiency. Therefore, when the refractive index of the cladding material changes, the hybrid mode would demonstrate quite different properties. Figure 5 shows the energy density distributions of a typical hybrid trench structure with $W_d =$ 300 nm and $W_s = 50$ nm. It is clear seen that, due to the high refractive index contrast and the relatively large dielectric core size, most of the mode energy can be stored inside the dielectric core when the refractive index of the cladding is small (e.g. see the results of $n_c = 1$ and $n_c = 1.5$). As the cladding's refractive index continuously gets larger, more energy would be released from the dielectric core into the low-index gap and the metal region (e.g. see $n_c =$ 2). Finally, when n_c gets very large (approaching n_d , i.e. 3.5), a large portion of the mode energy will be located around the sharp corners of the metal trench, leading to an increased modal loss.



Fig. 5. Energy density distribution of the hybrid trench waveguide ($W_d = 300$ nm) with different cladding material.

The confinement factor, defined as the power ratio inside each element to the total power of the waveguide, is a widely employed parameter to evaluate the confinement capability of the waveguide. In Fig. 6 we show the dependence of the confinement factor on n_c for various core size. It is illustrated that, for configurations with relatively large dielectric cores, a large portion of the mode energy can be resided in the dielectric core when n_c is small, which is consistent with the energy distributions shown in Fig. 5. This energy ratio would decrease monotonically when n_c gradually increases, as also can be seen from the distributions in Fig. 5. On the other hand, when the dielectric core size is relatively small (e.g. $W_d = 100$ nm), the energy inside the core is quite limited, and it demonstrate non-monotonic trends with the changing of the cladding's refractive index, which indicates the existence of an optimized n_c regarding the energy confined inside the dielectric core.



Fig. 6. The effect of the cladding layer's refractive index on the confinement factor for waveguides with different geometries: (a) confinement factor in the dielectric core; (b) confinement factor in the Ag trench; (c) confinement factor in the cladding and the gap region.

Figure 7 shows the relationship between the imaginary part of the modal effective index and the core width. Because the propagation length inversely proportional to the imaginary part of the waveguide's modal effective index, point A and point B show two critical values for W_d . Considering $W_s = 50$ nm, when $W_d > A$, the propagation length is decreasing monotonically with the increasing of the cladding's refractive index. On other hand, when W_d <A, the propagation length will increase first before it decreases with the increasing of n_c. It is also noted from Fig. 7 that, when $W_d < B$, the propagation length of $W_s = 50$ nm waveguide will always larger than that of $W_s = 25$ nm.



Fig. 7. Relationship between the imaginary part of modal effective index and the core width.

Finally, we show the dependence of propagation length and transmission loss on the height of the dielectric core. It is clearly seen in Fig. 8 that, increasing the height of the dielectric core could lead to the monotonic increase of the modal loss and thus decrease in the propagation length It is noted that, within the considered physical dimensions, low modal loss could be readily achieved, indicating reasonable propagation lengths for practical applications.



Fig. 8. Dependence of propagation length of the hybrid mode on the height of the trench (dielectric core).

The above discussions clearly reveal the nice optical performance of the proposed structure at the telecom wavelength. Compared to the metallic V-groove structure, the hybrid metal trench waveguide could realize much tighter confinement of light owing to the strong hybridization between the dielectric and plasmonic modes [31, 34]. On the other hand, the proposed hybrid waveguide enables much lower propagation loss than the standard metal slot waveguide [35], while maintaining subwavelength mode size. Here, it is also worth mentioning that, considering the practical fabrication processes, the proposed hybrid waveguide could have non-vertical metallic sidewalls (i.e. trapezoidal shaped metal trench). It is shown that similar modal behaviors can be obtained for these modifed hybrid trench waveguides, and their propagation losses can be even smaller than the rectangular-type structures studied here.

3. Conclusions

In summary, we have proposed and investigated a novel plasmonic waveguide that comprises a high-index dielectric core embedded inside a rectangular-shaped metallic trench structure. Numerical simulations reveal that the dielectric core can concentrate a large portion of the mode energy and also greatly reduce the propagation loss of the traditional metal trench waveguide. Through optimization of the dielectric core size and the refractive index of the cladding material, long propagation distance with strong optical confinement could be achieved simultaneously at the telecom wavelength. The studied hybrid structure could be employed to build highly efficient passive devices and may find other interesting applications as well.

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