

Spin-Orbit *Ab Initio* Investigation of the Photodissociation of Dibromomethane in the Gas and Solution Phases

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Abstract: A clear and reliable theoretical investigation on dibromomethane (CH_2Br_2) photodissociation is desired. The calculation must consider: (i) relativistic effects; (ii) the potential energy curves (PECs) of spin-orbit coupling states; (iii) geometry optimization by the method with both static and dynamic electron correlations; (iv) solvent effects on the photodissociation in the solution. All these have been considered in this study by state-of-the-art quantum chemical calculations. The experimentally observed photodissociation in the gas phase with products of spin-orbit-coupled states, $\text{Br}(^2\text{P}_{3/2})$ and $\text{Br}^*(^2\text{P}_{1/2})$, was assigned by multi-state second order multiconfigurational perturbation theory in conjunction with spin-orbit interaction through complete active space state interaction (MS-CASPT2/CASSI-SO) PECs. The mechanisms of the experimentally observed photodissociation and photoisomerization in solvent were elucidated by the MS-CASPT2/CASSI-SO method combined with polarized continuum model of the solvent.

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Key words: photodissociation; potential energy curve; spin-orbit coupling; solvent effect; MS-CASPT2

Introduction

The interest in the study of the photodissociations of polyhalomethane molecules has increased considerably in the last years. Molecules such as CH_2XY ($\text{X}, \text{Y} = \text{Cl}, \text{Br}, \text{and I}$) have been observed in the troposphere, and consequently considered as important sources of reactive halogens in the atmosphere.^{1–4} CH_2Br_2 and CH_2I_2 have been widely used as reagents for cyclopropanation reactions with alkenes.^{5–8} From a fundamental viewpoint, photodissociation studies involving polyhalomethanes are helpful for exploring fundamental photochemical reaction mechanisms. For CH_2Br_2 , there are rich experimental literatures on its ultraviolet absorption and photoelectron spectra,^{9–12} thermochemical data of photofragmentations,^{13–16} and photodissociation dynamics.^{17–22} The photodissociation of CH_2Br_2 has been investigated by all kinds of experimental techniques from different angles. Phillips and coworkers¹⁷ obtained A-band resonance Raman spectra of CH_2Br_2 in the gas and solution phases and nanosecond time-resolved resonance Raman spectra of CH_2Br_2 photoproducts. Lee et al.¹⁹ determined a simple C—Br bond fission of CH_2Br_2 at 248 nm by using product translational spectroscopy. Zhang and coworkers²¹ investigated the photodissociation dynamics of CH_2Br_2 and measured the branching ratios of $\text{Br}^*(^2\text{P}_{1/2})/\text{Br}(^2\text{P}_{3/2})$ at 234 and 267 nm. Wei et al.²² observed Br_2 molecular elimination in 248 nm photolysis of CH_2Br_2 by

using cavity ring-down absorption spectroscopy. One common conclusion of those experiments is that a simple C—Br bond cleavage is a main fast photodissociation channel of CH_2Br_2 in both gas and solution phases. In addition, CH_2Br_2 may turn a photodissociation into a photoisomerization in the solution phase.

Although there are rich experimental investigations, the detailed photodissociation process and mechanism of CH_2Br_2 are still not clear. It is in dire need of a clear reliable theoretical calculation in both gas and solution phases. The reliable calculation must consider: (i) relativistic effects since bromine is a heavy atom; (ii) the PECs of spin-orbit coupling states, which are very necessary to clearly describe the photodissociation channels leading to the products $\text{Br}(^2\text{P}_{3/2})$ and $\text{Br}^*(^2\text{P}_{1/2})$; (iii) geometry optimization by the method with both static and dynamic electron correlations, since the complete active space

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Table 1. The Geometries of CH₂Br₂, CH₂Br–Br and the TS Between them (Bond Lengths in Å and Bond Angles in Degrees).

Species	State	Method	R(C–H)	R(C–Br)	R(Br–Br)	∠HCH	∠HCB	∠BrCB	∠CBBr
CH ₂ Br ₂	S ₀	CASPT2	1.091	1.925		112.2	108.0	112.9	
		CASPT2 + PCM	1.092	1.927		112.4	108.0	112.7	
		QCISD ^a	1.086	1.945		111.6		114.2	
		Exp.	1.097, ^b 1.085 ^c	1.925, ^b 1.933 ^c		110.9, ^b 112.4 ^c		112.9, ^b 112.6 ^c	
		CASPT2	1.094	2.257		119.9	112.8	78.3	
		CASPT2 + PCM	1.095	2.265		120.2	112.8	77.9	
	S ₂	CASPT2	1.093	2.229		116.5	98.1	149.0	
CH ₂ Br–Br	T ₁	CASPT2	1.094	2.213		117.8	112.9	82.3	
		CASPT2	1.096	2.182		112.8	99.7	144.4	
		CASPT2	1.084	1.786	2.634	121.3	117.3	36.1	120.3
		CASPT2 + PCM	1.084	1.783	2.649	121.6	117.5	36.4	120.0
		B3LYP ^d	1.078	1.768	2.674		118.1		122.0
	S ₁	MP2 ^e	1.082	1.785	2.782	121.7	117.4		118.3
		CASPT2	1.081	1.853	4.232	123.6	117.4	84.2	70.0
		CASPT2 + PCM	1.082	1.855	4.250	123.9	117.2	84.7	69.6
		CASPT2	1.082	1.785, 3.013	2.898	123.6	99.0, 118.2	68.9	76.0, 35.1
		CASPT2							

^aRef. 34.^{b,c}Results based on microwave studies in Refs. 32 and 33.^dRef. 17.^eRef. 35.

self-consistent field (CASSCF) optimization on the weak C–Br bond could produce a very big deviation in the equilibrium geometry²³; (iv) solvent effects on the photodissociation in the solution. In the current theoretical investigation on the photochemistry of CH₂Br₂, we considered all the above aspects.

Methodology and Computational Details

The geometrical optimizations and vibrational frequency analyses were performed by using the complete active space second order perturbation (CASPT2)²⁴ method. The multi-state CASPT2 (MS-CASPT2)²⁵ was used to calculate the T_v values and oscillator strengths (f) of the spin-free states. The T_v and f values of the spin-coupled states were evaluated by the MS-CASPT2/CASSI-SO approach²⁶ in conjunction with atomic mean-field integral (AMFI) approximation.²⁷ The so-called IPEA-shift with a value of 0.25 was used in all the (MS-)CASPT2 calculations.²⁸ The C [He] and Br [Ar] cores were not correlated at the MS-CASPT2 level. The experimentally employed solvent cyclohexane was modeled via polarized continuum model (PCM).²⁹ The selected active space comprises 12 electrons in 10 orbitals for CH₂Br₂ and CH₂Br–Br. The relativistic basis sets of the atomic natural orbital type, ANO-RCC,³⁰ were used for all the above calculations. For the geometry optimizations, frequency analyses, and calculations of PECs, the H, C and Br bases were contracted to 2s1p, 3s2p1d, and 5s4p2d1f, respectively. For the calculations of the T_v and f values at equilibrium structures, a larger contraction of 3s2p1d for H, 4s3p2d1f for C and 6s5p3d2f1g for Br was used. All the calculations were performed using the MOL-CAS 6.4 quantum chemistry software.³¹

Results and Discussion

Geometries of CH₂Br₂ and CH₂Br–Br

The first three singlet and first two triplet states of CH₂Br₂, and the first two singlet states of CH₂Br–Br were located by the CASPT2 optimizations. The presently located geometries were compared with the previously optimized and available experimental ones in Table 1. For the schematic molecular structures, see Figure 1. The CASPT2 optimized bond length R(C–Br) and bond angle ∠BrCB of the CH₂Br₂ ground state (¹A₁) are nearly the same as the experimentally detected ones,^{32,33} respectively. The CASPT2 calculated harmonic vibrational frequencies of the CH₂Br₂ and CH₂Br–Br ground states were listed in Tables 2 and 3, respectively. The present results were compared with the previous QCISD³⁴ and DFT¹⁷ calculated and experimental values.^{35–37} As shown in Tables 2 and 3, the CASPT2 calculated vibrational frequencies were well or reasonably consistent with the corresponding experimental values.

We also optimized the geometries of the ground and first singlet excited states of CH₂Br₂ and CH₂Br–Br in cyclohexane solvent by the CASPT2 + PCM method. By the comparison in Table 1, the solvent has little effect on the geometries of CH₂Br₂. However, it does affect the weak Br–Br bond of CH₂Br–Br, which is elongated 0.015 and 0.018 Å in the solution for the ¹A' and ¹A'' states, respectively. The harmonic vibrational frequencies of CH₂Br₂ and CH₂Br–Br ground states were also computed by the CASPT2 + PCM method. As shown in Tables 2 and 3, the calculated results of vibrational frequencies by considering solvent effects are in better agreement with the experimentally observed values.

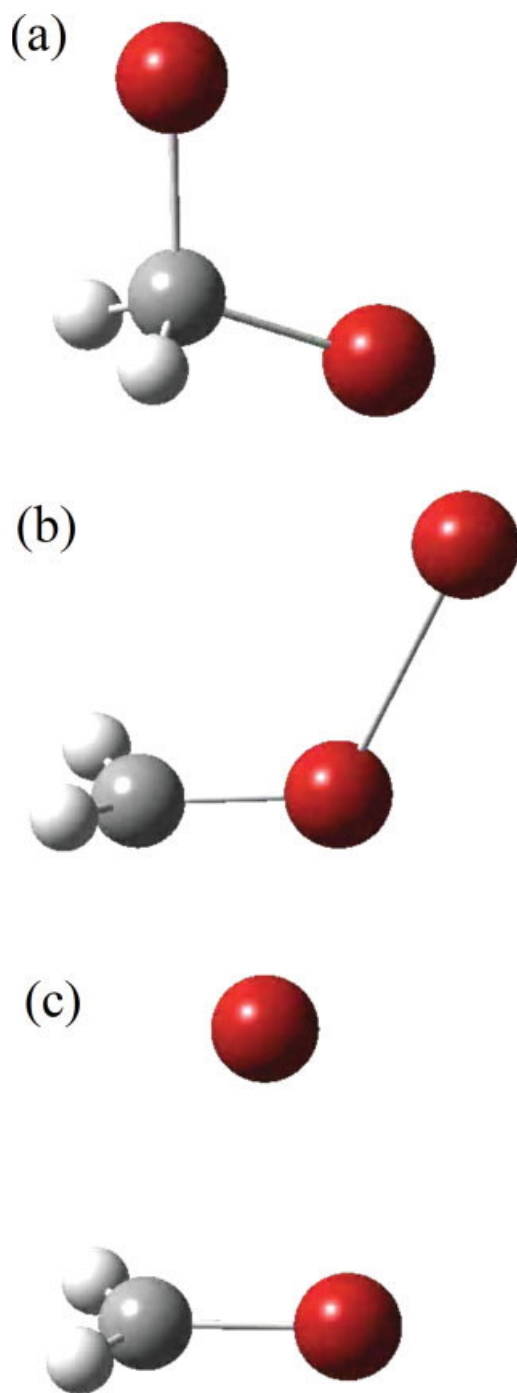


Figure 1. Schematic diagrams depicting the CASPT2 equilibrium geometries of (a) 1^1A_1 , 1^1B_1 , 1^1B_2 , 1^3B_1 , and 1^3B_2 of CH_2Br_2 ; (b) $1^1A'$ of $\text{CH}_2\text{Br}-\text{Br}$; (c) $1^1A'$ of $\text{CH}_2\text{Br}-\text{Br}$, this structure is similar to that of the TS between CH_2Br_2 and $\text{CH}_2\text{Br}-\text{Br}$.

Vertical Excitation Energies of CH_2Br_2

In 1975, Causley et al.⁹ observed the diffuse absorption bands in UV spectrum (denoted as A-bands) of CH_2Br_2 and assigned two band maxima $45,590\text{ cm}^{-1}$ (5.65 eV) and $51,490\text{ cm}^{-1}$ (6.38

Table 2. The Calculated CH_2Br_2 Harmonic Vibrational Frequencies in Comparison with the Experimental Values (in cm^{-1}).

Vibrational mode	CASPT2			Exp. ^b	Exp. ^c
	CASPT2	+ PCM	QCISD ^a		
CH sym. stretch	3118	3128	3020	3009	
CH_2 bend	1438	1434	1400	1382	1388 (1388)
CBr sym. stretch	589	569	562	588	587 (576)
CBr_2 bend	152	173	167	169	179 (174)
CH asym. stretch	3206	3230	3094	3073	
CH_2 rock	821	822	791	812	
CH_2 wag	1326	1227	1206	1195	
CBr asym. stretch	596	648	646	653	640 (637)
CH_2 twist	1121	1125	1094	1095	

^aScaled by 0.955.³⁴

^bRef. 36.

^cFrequencies observed in the gas phase and the ones in the parentheses observed in cyclohexane solution.¹⁷

eV) relevant to $n(b_1 + b_2) \rightarrow \sigma^*$ and $n(a_1 + a_2) \rightarrow \sigma^*$ transitions, respectively. Their experimental f values are 0.024 and 0.014,⁹ respectively. The T_v and f values of the first six singlet and the first six triplet excited states of CH_2Br_2 were calculated by the MS-CASPT2/CASPT2 method. The calculated results were compared with the previous TD-B3LYP^{21b} and MRDCI²² calculated ones in Table 4. Those very high states have nothing to do with our concerns and were not tabulated. The T_v (f) values of the first four singlet and the first four triplet excited states of CH_2Br_2 in cyclohexane were also calculated with consideration of solvent effects by the MS-CASPT2 + PCM/CASPT2 + PCM method (cf. Table 4). Compared with the results in the gas phase, the T_v values have not obvious changes in the solution phase. It is evidence that the solvent has little effect on the T_v values of CH_2Br_2 .

The T_v (f) values of seventeen spin-coupled states composed of the first five singlet and four triplet spin-free states in Table 4 were evaluated by the MS-CASPT2/CASSI-SO//CASPT2 and MS-CASPT2 + PCM/CASSI-SO//CASPT2 + PCM methods.

Table 3. The Calculated $\text{CH}_2\text{Br}-\text{Br}$ Harmonic Vibrational Frequencies in Comparison with the Experimental Values (in cm^{-1}).

Vibrational mode	CASPT2			Exp. ^a	Exp. ^b
	CASPT2	+ PCM	B3LYP ^a		
CH sym. stretch	3203	3204	3152		3030
CH_2 bend	1438	1409	1428		1334
CBr stretch	824	825	858		
CH_2 wag	670	697	738	690	684, 695
Br-Br stretch	193	190	180	176	
C-Br-Br bend	138	129	133	146	
CH asym. stretch	3358	3356	3286		3156
CH_2 rock	972	957	966	960	
CH_2 twist	446	430	468	480	

^aRef. 17.

^bFrom infrared absorption spectra in low-temperature matrixes in Refs. 35 and 37.

Table 4. The Calculated T_v (in eV) and f (in parentheses). Values of the First Four Singlet and the First Four Triplet Excited States of CH_2Br_2 .

State	MS-CASPT2// CASPT2	MS-CASPT2 + PCM// CASPT2 + PCM	TD-UB3LYP// UB3LYP ^a	MRDCI// B3LYP ^b
1^1A_1	0.00	0.00	0.00	0.00
1^1B_1	5.64 (2.69 E – 02)	5.63 (2.78 E – 02)	5.34 (0.0001)	5.94
1^1B_2	5.80 (6.98 E – 03)	5.81 (6.67 E – 03)	5.36 (0.0290)	5.96
1^1A_2	5.99 (0.00 E + 00)	5.98 (0.00 E + 00)	5.69 (0.0000)	
2^1A_1	6.29 (2.45 E – 04)	6.28 (1.25 E – 04)	5.90 (0.0013)	
1^3B_1	5.19	5.19	4.81	5.32
1^3B_2	5.26	5.25	4.94	5.48
1^3A_2	5.50	5.49	5.21	
1^3A_1	5.71	5.74	5.34	

^aRef. 21b.^bRef. 22.

The calculated T_v and f values and the compositions of the spin-free states in each spin-coupled state were listed in Table 5. According to the T_v and f values and the transition properties, we assign the experimentally⁹ observed two band maxima 45,590 and 51,490 cm^{-1} to 4B_1 and 5A_1 states, respectively.

Vertical Excitation Energies of $\text{CH}_2\text{Br}-\text{Br}$

We calculated the T_v (f) values of the first six singlet and the first six triplet excited states of $\text{CH}_2\text{Br}-\text{Br}$ by the MS-CASPT2//CASPT2 method. The values of the first three singlet and the first three triplet excited states were listed in Table 6

and compared with the previous URPA calculated results.¹⁷ We also calculated the T_v (f) values of the first three singlet and the first three triplet excited states of $\text{CH}_2\text{Br}-\text{Br}$ in cyclohexane by the MS-CASPT2 + PCM//CASPT2 + PCM method, see Table 6. Compared with the results in the gas phase, the T_v values have a decrease tendency in the solution phase.

The T_v (f) values of thirteen spin-coupled states composed of the first four singlet and three triplet spin-free states in Table 6 were evaluated by the MS-CASPT2/CASSI-SO//CASPT2 and MS-CASPT2 + PCM/CASSI-SO//CASPT2 + PCM methods. Table 7 listed the calculated T_v and f values and the compositions of the spin-free states in each spin-coupled state. As shown

Table 5. The Calculated T_v (in eV) and f (in parentheses) Values and Characteristics of the Low-Lying Spin-Coupled States of CH_2Br_2 in the Gas and Solution Phases.

State	MS-CASPT2/CASSI-SO//CASPT2		MS-CASPT2 + PCM/CASSI-SO//CASPT2 + PCM	
	$T_v(f)$	Characteristic	$T_v(f)$	Characteristic
1A_1	0.00	99.9% 1^1A_1	0.00	99.9% 1^1A_1
1B_2	5.24 (2.60 E – 06)	91.2% 1^3B_1 , 6.8% 1^3B_2 , 1.9% 1^3A_2	5.23 (3.29 E – 06)	91.8% 1^3B_1 , 6.4% 1^3B_2 , 1.8% 1^3A_2
1A_2	5.24 (0.00 E + 00)	92.8% 1^3B_1 , 6.4% 1^3B_2	5.24 (0.00 E + 00)	93.4% 1^3B_1 , 5.8% 1^3B_2
2A_1	5.24 (4.19 E – 05)	96.0% 1^3B_1 , 2.2% 1^3A_2	5.24 (3.89 E – 05)	96.1% 1^3B_1 , 2.0% 1^3A_2
1B_1	5.47 (7.46 E – 04)	92.5% 1^3B_2 , 6.0% 1^3A_1 , 1.5% 1^1B_1	5.48 (7.61 E – 04)	92.0% 1^3B_2 , 6.4% 1^3A_1 , 1.5% 1^1B_1
2A_2	5.49 (0.00 E + 00)	85.8% 1^3B_2 , 9.4% 1^3A_1 , 4.8% 1^3B_1	5.49 (0.00 E + 00)	85.6% 1^3B_2 , 10.0% 1^3A_1 , 4.2% 1^3B_1
3A_1	5.51 (5.85 E – 07)	91.0% 1^3B_2 , 5.8% 1^3B_1 , 1.7% 1^3A_2	5.51 (1.12 E – 06)	91.2% 1^3B_2 , 5.2% 1^3B_1 , 1.7% 1^3A_2 , 1.8% 2^1A_1
2B_1	5.65 (1.92 E – 03)	84.6% 1^3A_2 , 9.6% 1^1B_1 , 5.8% 1^3A_1	5.66 (2.12 E – 03)	83.6% 1^3A_2 , 9.8% 1^1B_1 , 6.4% 1^3A_1
4A_1	5.68 (4.30 E – 05)	92.8% 1^3A_2 , 3.8% 1^3A_1 , 2.4% 1^3B_1	5.69 (4.62 E – 05)	92.0% 1^3A_2 , 4.6% 1^3A_1 , 2.3% 1^3B_1
2B_2	5.69 (1.30 E – 05)	96.0% 1^3A_2 , 2.8% 1^3B_1 , 1.0% 1^3B_2	5.70 (1.41 E – 05)	96.1% 1^3A_2 , 2.6% 1^3B_1 , 1.0% 1^3B_2
3B_2	5.77 (8.28 E – 04)	80.0% 1^3A_1 , 15.1% 1^1B_2 , 4.6% 1^3A_2	5.77 (7.18 E – 04)	80.2% 1^3A_1 , 13.8% 1^1B_2 , 5.8% 1^3A_2
3B_1	5.82 (5.36 E – 03)	74.2% 1^3A_1 , 11.6% 1^3A_2 , 11.0% 1^1B_1 , 3.2% 1^3B_2	5.81 (5.11 E – 03)	74.2% 1^3A_1 , 12.4% 1^3A_2 , 9.8% 1^1B_1 , 3.6% 1^3B_2
3A_2	5.83 (0.00 E + 00)	89.7% 1^3A_1 , 7.8% 1^3B_2 , 2.0% 1^3B_1	5.83 (0.00 E + 00)	89.1% 1^3A_1 , 8.4% 1^3B_2 , 2.0% 1^3B_1
4B_1	5.94 (2.04 E – 02)	77.5% 1^1B_1 , 14.0% 1^3A_1 , 4.3% 1^3B_2 , 3.6% 1^3A_2	5.94 (2.17 E – 02)	78.4% 1^1B_1 , 12.8% 1^3A_1 , 4.3% 1^3B_2 , 3.8% 1^3A_2
4B_2	6.02 (4.73 E – 03)	83.0% 1^1B_2 , 15.0% 1^3A_1 , 1.5% 1^3B_1	6.03 (4.58 E – 03)	84.3% 1^1B_2 , 13.8% 1^3A_1 , 1.4% 1^3B_1
4A_2	6.23 (0.00 E + 00)	98.9% 1^1A_2	6.24 (0.00 E + 00)	98.8% 1^1A_2
5A_1	6.36 (6.80 E – 03)	98.1% 2^1A_1	6.27 (4.10 E – 03)	97.7% 2^1A_1 , 1.4% 1^1B_2

Table 6. The Calculated T_v (in eV) and f (in parentheses) Values of the First Three Singlet and the First Three Triplet Excited States of $\text{CH}_2\text{Br}-\text{Br}$.

State	MS-CASPT2// CASPT2	MS-CASPT2 + PCM// CASPT2 + PCM	URPA// UB3LYP ^a
$1^1\text{A}'$	0.00	0.00	0.00
$1^1\text{A}''$	2.20 (2.99 E - 06)	2.12 (1.69 E - 06)	3.18 (0.0002)
$2^1\text{A}'$	2.39 (3.65 E - 04)	2.32 (4.32 E - 04)	3.31 (0.3380)
$3^1\text{A}''$	4.00 (4.62 E - 01)	3.94 (5.09 E - 01)	3.48 (0.1919)
$1^3\text{A}'$	1.28	1.24	
$1^3\text{A}''$	2.16	2.07	
$2^3\text{A}'$	2.36	2.29	

^aRef. 17.

in Table 7, the predicted f value of $7\text{A}'$ is much larger than the others, which should be responsible for the experimentally observed absorption band ~ 360 nm (3.44 eV).¹⁷ However, the calculated T_v values are about 0.50–0.60 eV larger than the experimental value. Similar problems were encountered before for $\text{CH}_2\text{I}-\text{I}$,³⁸ $\text{CH}_2\text{Br}-\text{I}$,²³ and $\text{CH}_2\text{I}-\text{Br}$.²³ Detailed theoretical calculation is undergoing for a reliable quantitative explanation.

PECs of the Spin-Free and Spin-Coupled States of CH_2Br_2

The ground state CH_2Br_2 was constrainedly optimized at different C—Br bond distances using the CASPT2 method, and the energies of the first five singlet and the first four triplet states were calculated at the corresponding optimized ground state geometries by the MS-CASPT2 method. Thus we obtained the PECs of CH_2Br_2 including nine spin-free states, see Figure 2a. As shown in Figure 2a, the photodissociation products of higher excited states (1^1A_2 , 2^1A_1 , and 1^3A_1) are not related with the ground-state products, i.e., the experimentally observed fast dis-

sociation products ($\text{CH}_2\text{Br}\cdot + \text{Br}(^2\text{P}_{3/2})$ and $\text{CH}_2\text{Br}\cdot + \text{Br}^*(^2\text{P}_{1/2})$). So we do not need to consider these higher excited states for interpreting the photodissociation mechanism of CH_2Br_2 . The MS-CASPT2//CASPT2 calculated T_v values of the first and second singlet excited states to be 5.64 and 5.80 eV, respectively, as listed in Table 4. Obviously, the experimentally employed 234 nm (5.30 eV),²¹ 248 nm^{19,22} and 267 nm²¹ photon energy cannot reach 1^1B_1 and 1^1B_2 states vertically. The CASPT2 computed adiabatic excited energy (T_0) for 1^1B_1 and 1^1B_2 states of CH_2Br_2 is 4.90 and 4.91 eV, respectively. Therefore, the experimentally employed photon energy can make CH_2Br_2 transit from the ground state to the first two singlet excited states at most. As presented in Figure 2a, the PECs of the first two excited singlet and the first three triplet states are repulsive along the C—Br bond stretching coordinate and lead to the ground state products.

Figure 2a cannot give more information on the product of $\text{Br}^*(^2\text{P}_{1/2})$ or $\text{Br}(^2\text{P}_{3/2})$. Hence, the PECs of the corresponding seventeen spin-coupled states composed of the first five singlet and four triplet spin-free states in Table 4 were evaluated by the MS-CASPT2/CASSI-SO method. Those states higher than the fifteenth state have nothing to do with our concerns about the experimental observed photodissociation channels,^{19,21,22} so we just included the PECs of the first fifteen states in Figure 2b. The spin-orbit-coupled PECs were drawn diabatically according to the composition of the spin-free states in the spin-coupled states. As clearly shown in Figure 2b, there are three groups of dissociation products as the C—Br bond stretches to the dissociation limit. The highest group originating from 3B_2 , 3B_1 , and 3A_2 states (mainly composed of 1^3A_1) connects to very high excited states and is not related to the experimental observation. The energy difference between the lower two groups of dissociation products is 0.43 eV at $R(\text{C}-\text{Br}) = 5.0$ Å. This value agrees with the MS-CASPT2/CASSI-SO computed 0.43 eV²³ and the experimentally detected 0.46 eV^{19,39} energy gap between $\text{Br}^*(^2\text{P}_{1/2})$ and $\text{Br}(^2\text{P}_{3/2})$. So, we assigned the ground group of

Table 7. The Calculated T_v (in eV) and f (in parentheses) Values and Characteristics of the Low-Lying Spin-Coupled States of $\text{CH}_2\text{Br}-\text{Br}$ in the Gas and Solution Phases.

State	MS-CASPT2/CASSI-SO//CASPT2		State	MS-CASPT2 + PCM/CASSI-SO//CASPT2 + PCM	
	$T_v(f)$	Characteristic		$T_v(f)$	Characteristic
$1\text{A}'$	0.00	99.6% $1^1\text{A}'$	$1\text{A}'$	0.00	99.6% $1^1\text{A}'$
$1\text{A}''$	1.28 (0.00 E + 00)	97.9% $1^3\text{A}'$	$2\text{A}'$	1.29 (6.20 E - 06)	98.0% $1^3\text{A}'$
$2\text{A}'$	1.28 (1.22 E - 05)	97.8% $1^3\text{A}'$	$1\text{A}''$	1.29 (0.00 E + 00)	98.1% $1^3\text{A}'$
$2\text{A}''$	1.28 (3.05 E - 08)	98.0% $1^3\text{A}'$	$2\text{A}''$	1.29 (1.01 E - 08)	97.9% $1^3\text{A}'$
$3\text{A}'$	2.15 (3.86 E - 04)	79.2% $1^3\text{A}''$, 19.4% $2^3\text{A}'$	$3\text{A}'$	2.19 (4.26 E - 04)	79.0% $1^3\text{A}''$, 18.4% $2^3\text{A}'$
$3\text{A}''$	2.15 (0.00 E + 00)	76.8% $1^3\text{A}''$, 23.0% $2^3\text{A}'$	$3\text{A}''$	2.19 (0.00 E + 00)	76.4% $1^3\text{A}''$, 23.3% $2^3\text{A}'$
$4\text{A}'$	2.16 (2.46 E - 05)	79.7% $1^3\text{A}''$, 18.4% $2^3\text{A}'$	$4\text{A}'$	2.20 (5.28 E - 05)	78.2% $1^3\text{A}''$, 19.0% $2^3\text{A}'$
$4\text{A}''$	2.17 (2.12 E - 07)	74.1% $1^1\text{A}''$, 25.8% $2^3\text{A}'$	$4\text{A}''$	2.20 (6.38 E - 07)	75.1% $1^1\text{A}''$, 24.4% $2^3\text{A}'$
$5\text{A}'$	2.47 (4.31 E - 03)	70.8% $2^3\text{A}'$, 20.5% $1^3\text{A}''$, 7.9% $2^1\text{A}'$	$5\text{A}'$	2.51 (4.80 E - 03)	66.6% $2^3\text{A}'$, 20.7% $1^3\text{A}''$, 11.6% $2^1\text{A}'$
$5\text{A}''$	2.50 (1.50 E - 08)	75.9% $2^3\text{A}'$, 22.0% $1^3\text{A}''$	$5\text{A}''$	2.53 (2.07 E - 08)	75.9% $2^3\text{A}'$, 22.2% $1^3\text{A}''$
$6\text{A}''$	2.50 (3.00 E - 08)	73.3% $2^3\text{A}'$, 24.7% $1^1\text{A}''$	$6\text{A}''$	2.54 (3.87 E - 07)	74.5% $2^3\text{A}'$, 23.5% $1^1\text{A}''$
$6\text{A}'$	2.52 (5.60 E - 05)	71.4% $2^1\text{A}'$, 19.0% $1^3\text{A}''$, 8.0% $2^3\text{A}'$	$6\text{A}'$	2.55 (1.85 E - 04)	65.9% $2^1\text{A}'$, 20.7% $1^3\text{A}''$, 12.0% $2^3\text{A}'$
$7\text{A}'$	4.05 (6.17 E - 01)	99.5% $3^1\text{A}'$	$7\text{A}'$	3.95 (5.43 E - 01)	99.3% $3^1\text{A}'$

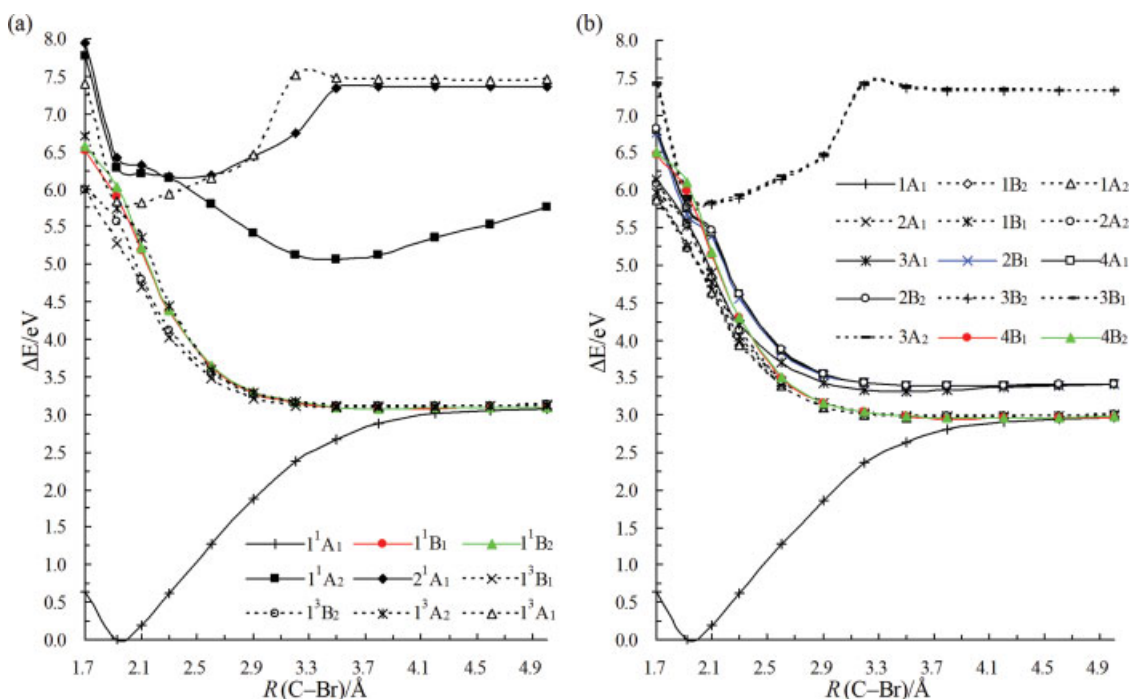


Figure 2. (a) The spin-free MS-CASPT2/CASPT2 PECs and (b) the spin-orbit-coupled MS-CASPT2/CASSI-SO/CASPT2 PECs, with respect to the C—Br bond coordinate of CH_2Br_2 .

dissociation products to $\text{CH}_2\text{Br}\cdot + \text{Br}^2(\text{P}_{3/2})$, and the upper group to $\text{CH}_2\text{Br}\cdot + \text{Br}^*(^2\text{P}_{1/2})$. As shown in Figure 2b, products Br are originally from the 1B_2 , 1A_2 , 2A_1 , 1B_1 , 2A_2 , 4B_1 , and 4B_2 states, while products Br^* are from the 3A_1 , 2B_1 , 4A_1 , and 2B_2 states. Many of these states have absolute zero or very small f values, which have no or very little contribution to the observed dissociation channels. Among them, the 2B_1 , 4B_1 , and 4B_2 states are of f values larger than $1.0\text{E}-03$. These three states should be important for assigning the experimentally observed dissociation channels. According to Table 5, most states have strong spin-orbit coupling. States 2B_1 and 4B_1 are of 9.6% and 77.5% composition of 1^1B_1 , respectively. State 4B_2 is mainly composed of the 1^1B_2 state. As presented in Figure 2b, state 2B_1 directly dissociates to $\text{CH}_2\text{Br}\cdot + \text{Br}^*(^2\text{P}_{1/2})$, while states 4B_1 and 4B_2 dissociate to $\text{CH}_2\text{Br}\cdot + \text{Br}^2(\text{P}_{3/2})$. As shown in Table 5, states 4B_2 and 2B_1 have f values at the same level, which are much smaller than the f value of state 4B_1 . So, the photodissociation with Br recoil mostly occurs on the B_1 state, and the yield of Br is larger than that of Br^* . This conclusion drawn from the present calculations agrees with the experimental observations.^{17,19,21,22} The experiment¹⁹ determined anisotropy parameter β value to be 1.0 ± 0.2 , which indicates that the photodissociation with Br recoil mostly occurs on the B_1 state.

To consider solvent effects, the spin-free and spin-coupled states' PECs of CH_2Br_2 were also recalculated by the MS-CASPT2 + PCM and MS-CASPT2 + PCM/CASSI-SO methods, respectively. As shown in Figure S1 (see Supplementary Material Section), the PECs of CH_2Br_2 with consideration of solvent effects are very similar to the ones in the gas phase.

Consequently, solvent effects do not affect the dissociation mechanism of CH_2Br_2 .

Isomerization between CH_2Br_2 and $\text{CH}_2\text{Br}-\text{Br}$

Phillips and coworkers¹⁷ reported solvent effects on the A-band photodissociation of CH_2Br_2 by observing resonance Raman spectra in the solution phase. They considered that CH_2Br_2 may turn a photodissociation into a photoisomerization in the solution

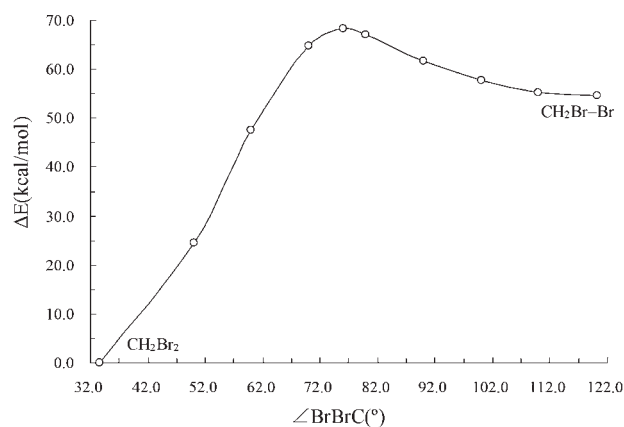


Figure 3. The CASPT2 + PCM/CASPT2 calculated PEC diagram depicting the isomerization between CH_2Br_2 and $\text{CH}_2\text{Br}-\text{Br}$ in the solution phase.

phase. In the isomerization reaction, it is evident that the isomer $\text{CH}_2\text{Br}-\text{Br}$ is responsible for the ~ 360 nm transient absorption band observed upon ultraviolet photolysis of CH_2Br_2 .¹⁷ This absorption was assigned to the $7\text{A}'$ state of $\text{CH}_2\text{Br}-\text{Br}$ in section 3.3. This process cannot occur on the excited states, since the first several excited states are repulsive along the C—Br bond of CH_2Br_2 (see Fig. 2a). When the C—Br bond of CH_2Br_2 stretches to $\text{CH}_2\text{Br}\cdot + \text{Br}$, the solvent cage effect induces the recombination of the $\text{CH}_2\text{Br}\cdot$ and Br fragments to $\text{CH}_2\text{Br}-\text{Br}$. We chose the bond angle $\angle\text{BrBrC}$ as reaction coordinate and calculated the isomerization reaction path by the CASPT2 + PCM/CASPT2 method. As presented in Figure 3, the CASPT2 + PCM calculated activation energy from $\text{CH}_2\text{Br}-\text{Br}$ to CH_2Br_2 is 13.7 kcal/mol, and 68.3 kcal/mol for the inverse reaction. A transition state (TS) for the isomerization process was located by the CASPT2 full optimization (see Fig. 1 and Table 1) with only one imaginary frequency of 587 cm^{-1} .

Conclusions

The photodissociation of CH_2Br_2 in the gas and solution phases has been investigated by the MS-CASPT2/CASSI-SO and MS-CASPT2 + PCM/CASSI-SO methods, respectively. The calculated PECs of the spin-coupled states clearly assign the experimentally observed photodissociation channels leading to $\text{CH}_2\text{Br} + \text{Br}({}^2\text{P}_{3/2})$ and $\text{CH}_2\text{Br}\cdot + \text{Br}({}^2\text{P}_{1/2})$ in both gas and solution phases. For the photolysis of CH_2Br_2 in the solution, there is also an isomerization from CH_2Br_2 to $\text{CH}_2\text{Br}-\text{Br}$. But $\text{CH}_2\text{Br}-\text{Br}$ is much more unstable and prefers to turn back to CH_2Br_2 .

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